

Chapter 1

1.1 Consider a narrow band acoustic spectrum expressed in terms of angular frequency $G_{pp}(\omega)$ where the mean square fluctuation is divided up by its contribution to 1 rad/s intervals (i.e. the spectral density per rad/s). We can define the sound pressure level in decibels for this spectral density as $10 \log_{10} \left(\frac{G_{pp}(\omega)}{p_{ref}^2 / (\text{rad/s})} \right)$ where, for air $p_{ref}^2 = 4 \times 10^{-10} (\text{Pa})^2$. If $G_{pp}(\omega)$ varies with frequency as B/ω^2 , and $B = 25 (\text{Pa})^2/\text{s}$, what is the overall A-weighted sound pressure level? *Hint: you will need to research the formula for A-weighting.* [Worked example solution](#)

1.2 A noise source has a narrowband sound spectrum (1 Hz intervals) which varies as $-20 \log_{10}(f) + 140$ where f is in Hz. This equation represents the spectral density per Hz in decibels ($\text{dB}(\text{Hz}) = 10 \log_{10}((p_{rms}^2/\text{Hz})/(p_{ref}^2/\text{Hz}))$) with a reference pressure of $20 \mu\text{Pa}$. This equation is defined only in the range 100 Hz to 10 kHz. Assume that sound is not produced outside of this range. Compute and plot the $1/3^{\text{rd}}$ octave-band spectra for this narrowband spectrum and the overall sound pressure level in decibels.

1.3. A jet engine produces tones (single frequency sound waves) at integer multiples of 200Hz. Assuming the tones have equal amplitude; (a) plot the $1/3^{\text{rd}}$ octave SPL spectrum from 200Hz to 20kHz, (b) plot the A-weighted $1/3^{\text{rd}}$ octave spectrum, (c) calculate the overall SPL and the overall A-weighted SPL. Your plots should use a logarithmic frequency scale. Explain why $1/3^{\text{rd}}$ octave band spectra are unlikely to be useful in general for representing tone-noise sources. *Hint: you will need to research the formula for A-weighting.*

Solution Problem 1.1

$$G_{pp}(\omega) = B/\omega^2$$

A-weighting (see <https://en.wikipedia.org/wiki/A-weighting>) is

$$A(f) = 20 \log_{10} R_A(f) + 2.00$$

or

$$A(f) = 10 \log_{10}(1.259^2 R_A(f)^2)$$

where

$$R_A(f) = \frac{12194^2 f^4}{(f^2 + 20.6^2)(f^2 + 12194^2)\sqrt{(f^2 + 107.7^2)(f^2 + 737.9^2)}}$$

Now, we have per (rad/s) that

$$SPL = 10 \log_{10} \left(\frac{G_{pp}(\omega)}{4 \times 10^{-10}} \right)$$

and

$$\begin{aligned} SPL_A &= 10 \log_{10} \left(\frac{G_{pp}(\omega)}{4 \times 10^{-10}} \right) + 10 \log_{10}(1.259^2 R_A(f)^2) \\ &= 10 \log_{10} \left(\frac{1.259^2 R_A(f)^2 G_{pp}(\omega)}{4 \times 10^{-10}} \right) \end{aligned}$$

where $f = 2\pi\omega$. Thus,

$$\begin{aligned} OASPL_A &= 10 \log_{10} \left[\int_0^{\infty} \frac{1.259^2 R_A(f)^2 G_{pp}(\omega)}{4 \times 10^{-10}} d\omega \right] \\ &= 10 \log_{10} \left[\frac{1.259^2}{4 \times 10^{-10}} \int_0^{\infty} \frac{12194^4 f^8 \times 2\pi \times 25 / (2\pi f)^2}{(f^2 + 20.6^2)^2 (f^2 + 12194^2)^2 (f^2 + 107.7^2)(f^2 + 737.9^2)} df \right] \end{aligned}$$

Since, $G_{pp}(\omega)d\omega = G_{pp}(f)df$, so

$$OASPL_A = 10 \log_{10} \left[\frac{12194^4 \times 1.259^2 \times 25}{8\pi \times 10^{-10}} \int_0^{\infty} \frac{f^6}{(f^2 + 20.6^2)^2 (f^2 + 12194^2)^2 (f^2 + 107.7^2)(f^2 + 737.9^2)} df \right]$$

Using Wolfram Alpha, the integral is evaluated as 7.52482×10^{-20} , so

$$OASPL_A = 10 \log_{10} \left[\frac{12194^4 \times 1.259^2 \times 25}{8\pi \times 10^{-10}} \times 7.52482 \times 10^{-20} \right] = 74.2 \text{ dB(A)}$$

Note that A-weighting is only defined for 20Hz to 20 kHz, but contributions outside this range to the above integral are negligible. Also note that a conventional approach (producing almost the same answer) would be to first calculate the 1/3rd octave band SPL from G_{pp} , then A-weight these values at the band center frequencies, and then integrate the mean-square pressures implied by these weighted values to get the final OASPLA.